















These lectures are mainly aimed at understanding primary conversion of wave-energy









A particular waterparticle moves once around its
circle in a time T (the "period").
The troughs are wider than the crests.
The difference is negligible for small waves

$$a \ll L/2\pi = 1/k$$
.
Then the water surface has a sinusoidal shape.
"Small waves"= "linear waves" (linear theory applicable)
 $\lambda = L$ is the "wavelength"
 $k = 2\pi/L$ is the "angular repetency" (or "uluve number")





Balance of flow:
$$wb2h = wal/\pi$$
 $\left[\frac{a}{b} = kh = \frac{2\pi h}{L}\right]$
On very shallow water the maximum values
of the water particle excursion and the water
velocity have a ratio
 $4/b = wa/wb = kh = 2\pi h/L$
between the vertical and horizontal components.
("shallow water" means $kh \ll 1$)
On deep water $(kh >> 1)$
• cincular orbits $b = a$
 $(V_x)_{max} = (V_z)_{max} = wa$





$$\dot{s}_{z} = V_{z} = -\omega a \sin(\omega t - kx)$$

$$\dot{s}_{z} = V_{x} = \omega b \cos(\omega t - kx)$$

$$\frac{horizontal acceleration}{gravity acceleration (g)} = tangent of angle of tilt$$
Maximum horizontal acceleration: $(\dot{v}_{x})_{max} = g \tan(ka) \approx g ka$
From orbit motion: $(\dot{v}_{x})_{max} = (\ddot{s}_{x})_{max} = \omega^{2}b$

$$\underbrace{\delta a = 0}_{fa} \qquad g ka = \omega^{2}b \quad \underbrace{\delta \omega^{2} = g k \frac{a}{b}}_{fa}$$
Phase velocity: $c = \frac{L}{T} = \frac{\omega}{L} = \frac{q}{\omega} (a/b) = \sqrt{\frac{q}{L}} (a/b)$
The wavelength: $L = 2\pi/k = \frac{q}{2\pi} T^{2} (a/b)$
The period: $T = 2\pi/\omega = \{2\pi, \frac{L}{g}, (b/a)\}^{1/2}$

For "deep" water (that is
$$kh \gg 1$$
):
we have $b/a = 1$ (circular orbit)
Then $c = \frac{q}{\omega} = \left\{\frac{q}{4k}\right\}^{1/2} = \frac{q}{2\pi}T = \left\{\frac{q}{2\pi}L\right\}^{1/2}$
 $L = cT = \frac{q}{2\pi}T^{2}$
 $\frac{q}{2\pi} = \frac{q.81 \text{ m/s}^{2}}{2\pi} = 1.56 \text{ m/s}^{2}$
Example: $T = 10s$; $c = 15.6 \text{ m/s}$. $L = 156 \text{ m}$
 $(= 56.2 \text{ km/h})$
The deep-water approximation is usually acceptable
if $h > L/3$ (kh > 2)

For shallow water (that is
$$kh \ll 1$$
);
we found $a/b = kh = 2\pi h/L$
 $co^2 = gk(a/b) = gkkh = ghk^2$
 $c = \frac{L}{T} = \frac{C}{K} = \sqrt{gh}$
 $L = cT = T\sqrt{gh}$
Example: $h = 6 \text{ m}$ $c = \sqrt{9.81 \cdot 6j0} = 7.7 \text{ m/s}$
 $(= 27.6 \text{ km/h})$
H $T=20s = L = 153 \text{ m}$
The shallow-water approximation is usually
acceptable if $h \leq L/20$ ($kh \leq 0.3$)

Superposition of two waves of the same

$$L H$$
amplitude but slightly different frequencies:

$$\gamma_{1} = a \cos(\omega_{1}t - k_{1}x) \qquad \gamma_{2} = a \cos(\omega_{2}t - k_{2}x)$$

$$\gamma_{(x,t)} = \gamma_{1}(x,t) + \gamma_{2}(x,t)$$
Recalling the trigonometric identity

$$\cos \alpha_{1} + \cos \alpha_{2} = 2 \cos \frac{\alpha_{1} - \alpha_{2}}{2} \cos \frac{\alpha_{1} + \alpha_{2}}{2}$$

$$\gamma_{2} = 2a \cos \left(\frac{\omega_{2} - \omega_{1}}{2}t - \frac{k_{2} - k_{1}}{2}x\right) \cos \left(\frac{\omega_{2} + \omega_{3}}{2}t - \frac{k_{2} + k_{2}}{2}x\right)$$

$$\gamma_{1} = a \cos(\omega_{1}t - k_{1}x) \qquad \gamma_{2} = a \cos(\omega_{2}t - k_{2}x)$$

$$\gamma(x,t) = \gamma_{1}(x,t) + \gamma_{2}(x,t)$$

$$\gamma = 2a \cos\left(\frac{\omega_{2} - \omega_{1}}{2}t - \frac{k_{2} - k_{1}}{2}x\right) \cos\left(\frac{\omega_{2} + \omega_{1}}{2}t - \frac{k_{2} + k_{2}}{2}x\right)$$
Set $\omega_{1} = \omega - \Delta \omega$ $\omega_{2} = \omega + \Delta \omega$

$$k_{1} = k - \Delta k \qquad k_{2} = k + \Delta k$$

$$\gamma = 2a \cos\left(a\omega t - \Delta k x\right) \cos\left(\omega t - kx\right)$$
Varies slowly
if $\Delta \omega \ll \omega$ (and hence $\Delta k \ll k$)
Resulting wave of "angular trequency" ω
and a slowly varying "amplitude"
 $2a \cos\left(a\omega t - \Delta k x\right)$
This "amplitude" propagates with a
speed
 $C_{g} = \frac{\Delta \omega}{\Delta k}$
which (it $\Delta \omega \gg 0$) is termed the group velocity

$$\frac{\text{horizontal acceleration}}{\text{gravity acceleration }(g)} = \text{tangent of angle of tilt}$$

$$\text{Maximum horizontal acceleration: } (\overset{}{V_{x}})_{max} = g \tan(ka) \approx g ka$$
From orbit motion: $(\overset{}{V_{x}})_{max} = (\overset{}{S_{x}})_{max} = \omega^{2}b$

$$\overbrace{=}^{a_{1}} & g ka = \omega^{2}b \quad \underbrace{\{\omega^{2} = gk, \frac{a}{b}\}}_{a}$$

$$Phase velocity: c = \frac{L}{T} = \frac{\omega}{J_{k}} = \frac{g}{\omega}(a/b) = \sqrt{\frac{g}{J_{k}}}(a/b)$$
The wavelength: $L = 2\pi/k = \frac{g}{2\pi}T^{2}(a/b)$
The period: $T = 2\pi/\omega = \left\{2\pi \frac{L}{g}(b/a)\right\}^{1/2}$

We found $w^2 = gk(a/b)$ (the "dispersion" equation) On deep water (a/b = 1) $w^2 = gk \Rightarrow 2wdw = gdk$ Phase velocity $c' = \frac{w}{k} = \frac{g}{w} = \sqrt{g/k}$ Group velocity $c_g = \frac{dw}{dk} = \frac{g}{2w} = \frac{1}{2}c$ Thus we have the important result that on deep water the group velocity is half of the phase velocity On very shallow water (a/b = kh) $w = k\sqrt{gh}$ $C_g = \frac{dw}{dk} = \frac{w}{k} = \sqrt{gh} = c$ The group velocity and phase velocity are equal and independent of the frequency so long as the water may be considered to be shallow $(w\sqrt{h/g} = kh <<1)$





Example: Assume that the sea is calm. Then, suddenly a storm develops 1=300 km from land. How long time afterwards can we record swells of pariod T= 14s at the shore ? What if the period is T= 10s? Assume that the water depth is more than 200 m.

Solution:
Deep-water formules are applicable because

$$L = 1.56 T^2 = 1.56 \cdot 14^2 = 306 \text{ m}$$
, that is
 $h > 200 \text{ m} > L/3 = 102 \text{ m}$
The group velocity is
 $C_g = \frac{1}{2}C = \frac{1}{2}\frac{L}{T} = \frac{1}{2}1.56 \times T = 0.78 T$
 $= 0.78 \times 14 = 10.9 \text{ m/s}$
Time before swell record
 $\Delta t = \frac{L}{C_g} = \frac{R}{0.78T} = \frac{300 \cdot 10^2}{0.78T} = \frac{3.8 \cdot 10^5}{T}$
 $T = 14 \text{ A}$: $\Delta t = \frac{3.8 \cdot 10^5}{14} = 27 \cdot 10^3 \text{ c} = 7.6 \text{ hours}$
 $T = 10 \text{ s}$ $\Delta t = \frac{3.8 \cdot 10^5}{40} = 38 \cdot 10^3 \text{ s} = 10.7 \text{ hours}$







For the case of deep water:
$$c_q = c_f/2 = \frac{q}{2\omega} = \frac{q}{4\pi}$$

The «wave-power level»:
 $\int = \frac{q q^2}{4\omega} |\eta|_{max}^2 = \frac{q q^2 T}{8\pi} |\eta|_{max}^2 \frac{q q^2 T}{32\pi} H^2 = (q 76 \frac{W}{3.000}) T H^2$
The wave height: $H = 2|\eta|_{max}$
Plane wave propagating on deep $\eta = A \cos(\omega t - k_X + \alpha)$
 $\int = \frac{q q^2}{4\omega} A^2 = \frac{q q^2}{8\pi f} A^2$
 $E = 2E_p = 2E_k = \frac{1}{2}ggA^2 = \frac{q q^2}{2}$
Multi-frequency $\eta = \sum_m A_m \cos(\omega_m t - k_m X + \alpha_m)$
 $sea wave: \int = \sum_m \frac{q q^2}{8\pi f_m} A_m^2$
 $E = \frac{1}{2}gg \sum_m A_m^2$

Multi-frequency
sea wave:
$$\gamma = \sum_{m} A_{m} \cos(\omega_{m} t - k_{m} x + d_{m})$$

 $\int = \sum_{m} \frac{Qq^{2}}{8\pi f_{m}} A_{m}^{2}$ $E = \frac{1}{2}gq \sum_{m} A_{m}^{2}$
More
general
sea wave: $E = qq \frac{\gamma}{2}(x, y, t) = gq \int S(f) df$
where we have introduced the real sea wave's
«energy spectrum» $S(f)$, for which the SI unit is
 m^{2}/Hz . The overbar denotes time average.
Spectrally defined «significant wave height»:
 $H_{mo} = 4\sqrt{m_{o}}$ $m_{o} = \int_{0}^{\infty} S(f) df$

$$\begin{split} \eta &= \sum_{m} A_{m} \cos \left(\omega_{m} t - k_{m} \times t \cdot d_{m} \right) \\ J &= \sum_{m} \frac{Q q^{2}}{8 \pi f_{m}} A_{m}^{2} \qquad E = \frac{1}{2} g q \sum_{m} A_{m}^{2} \\ \eta &= \eta \left(x_{1} \eta_{1} t \right) \\ E &= Q q \overline{\eta^{2}(x_{1} \eta_{1} t)} \equiv Q q \int_{0}^{\infty} S(q) df \end{split}$$
Spectrally defined «significant wave height»:
$$\begin{aligned} H_{mo} &= 4 \sqrt{m_{o}} \qquad m_{o} \equiv \int_{0}^{\infty} S(q) df \end{aligned}$$
Spectrally defined «wave-power level»:
$$\begin{aligned} J &= \frac{Q q^{2}}{4 \pi} \int_{0}^{\infty} \frac{S(q)}{q} df = \frac{Q q^{2}}{4 \pi} m_{q} \\ m_{-1} &= \int_{0}^{\infty} q^{-1} S(q) df \end{aligned}$$

Spectrally defined «significant wave height»:

$$H_{m_0} = 4 \sqrt{m_0} \qquad m_0 = \int S(f) df$$
Spectrally defined «wave-power level»:

$$J = \frac{99^2}{4\pi} \int \frac{S(f)}{f} df = \frac{99^2}{4\pi} m_1$$

$$m_{-1} = \int f^{-1} S(f) df$$
Spectral moment of order *j*: $m_{j} = \int f^{-j} S(f) df$
Spectrally defined
«energy period»: $T_J = T_{-1,0} = \frac{m_{-1}}{m_0}$
and «wave-
power level»: $J = \frac{99^2}{4\pi} T_J m_0 = \frac{99^2}{64\pi} T_J H_{m_0}^2$

Wind waves and swells

•Waves generated by wind are called *wind waves*. When the waves propagate outside their region of generation, they are called *swells* [in Norwegian: *dønning*]. Where the water is deep, swells can travel very large distances, for instance across oceans, almost without loss of energy.















<text><text><equation-block><text>

$$E = \rho g \int_{0}^{\infty} S(f) df \equiv \rho g H_{s}^{2} / 16$$

$$\gamma = \sum_{m} A_{m} \cos(\omega_{m} t - k_{m} x + d_{m})$$

$$J = \sum_{m} \frac{\varphi q^{2}}{8\pi f_{m}} A_{m}^{2} \qquad E = \frac{1}{2} g q \sum_{m} A_{m}^{2}$$

$$\gamma = \gamma (x, y, t)$$

$$E = \varphi q \qquad \overline{\gamma^{2}(x, y, t)} = \varphi q \int_{0}^{\infty} S(t) df$$
Spectrally defined «significant wave height»:
$$H_{m0} = 4 \sqrt{m_{0}} \qquad m_{0} = \int_{0}^{\infty} S(t) df$$
Spectrally defined «wave-power level»:
$$J = \frac{\varphi q^{2}}{4\pi} \int_{0}^{\infty} \frac{S(t)}{t} df = \frac{\varphi q^{2}}{4\pi} m_{-1}$$

$$m_{-1} = \int_{0}^{\infty} f^{-1} S(t) df$$





































For mathematical simplicity the damping force
is assumed proportional to the velocity
$$\dot{x} = \frac{dx}{dt}$$

 $R = "mechanical resistance"$
The spring force is assumed proportional to the
displacement x from equilibrium.
Newton: $m\ddot{x} = F + F_s + F_R = F - Sx - R\dot{x}$
 $m\ddot{x} + R\dot{x} + S\dot{x} = F$ (2
(External force balanced against 1) inertial force
amping force and 3) spring force)

$$m\ddot{x} + R\dot{x} + S\dot{x} = F$$
(2
(External force balanced against ¹) inertial force
²) damping force and ³) spring force)
Energy = Force × Displacement
Power = Force × Velocity
Power applied to the mechanical system

$$P = F\dot{x} = m\ddot{x}\dot{x} + R\dot{x}^{2} + S\dot{x}\dot{x} =$$

$$= m\dot{x}\frac{d\dot{x}}{dt} + R\dot{x}^{2} + S\frac{dx}{dt} \times =$$

$$= \frac{d}{dt}(\frac{1}{2}m\dot{x}^{2}) + R\dot{x}^{2} + \frac{d}{dt}(\frac{1}{2}Sx^{2})$$

$$= R\dot{x}^{2} + \frac{d}{dt}(W_{k} + W_{p})$$

Power applied to the mechanical system

$$P = F \dot{x} = m \ddot{x} \dot{x} + R \dot{x}^{2} + S \dot{x} \dot{x} =$$

$$= m \dot{x} \frac{d\dot{x}}{dt} + R \dot{x}^{2} + S \frac{dx}{dt} x =$$

$$= \frac{d\dot{y}}{dt} \left(\frac{1}{2}m \dot{x}^{2}\right) + R \dot{x}^{2} + \frac{d}{dt} \left(\frac{1}{2}S x^{2}\right)$$

$$= R \dot{x}^{2} + \frac{d}{dt} \left(W_{k} + W_{p}\right)$$

$$W_{k} = \frac{1}{2}m \dot{x}^{2} = kinetic energy$$

$$W_{p} = \frac{1}{4}S x^{2} = potential energy (of the spring)$$



$$P(t) = P_{R}(t) + (P_{k}(t) + P_{p}(t))$$

$$P_{R}(t) = -F_{R}(t) u(t) = R u^{2}$$

$$P_{R}(t) = -F_{R}(t) u(t) = m \dot{u} u = \frac{d}{dt} W_{k}(t)$$

$$P_{R}(t) = \frac{1}{2} m (u(t))^{2} - kinetic energy.$$

$$P_{p}(t) = -F_{s}(t) u(t) = S \times \dot{x} = \frac{d}{dt} W_{p}(t) \quad (t.32)$$

$$W_{p}(t) = \frac{1}{2} S (x(t))^{4} - potential energy.$$
Energy stored in the oscillating system:

$$W(t) = W_{k}(t) + W_{p}(t) \quad (2.35)$$

$$P_{k}(t) + P_{p}(t) = \frac{d}{dt} W(t) \quad (2.76)$$



콭 distance & from its aquilibrium position, the upward buoyancy force is reduced by SgAx. Thus there will be a restoring face $F_s = -ggA_w \times$ Excling in a direction to restore the equilibrium (Note: position and force are chosen to be positive upwards) F = -SxHydrostatic stiffness (buoyancy stiffness) S= 29 Aw If the floating body is axisymmetric and of radius a (diameter 2a) the "water plane area" is Aw = Ta" 5= 89 17 22











In this simple example, at optimum radiated-wave generation, the maximum absorbed energy equals 100 percent of the incident wave energy. Note also that the required, optimum, radiated wave has the same amplitude as the incident wave. Thus,



Observe that, in order to absorb, from the sea, the theoretically maximum wave power, it is necessary that the waveabsorbing oscillating system, at optimum, has an *ability* to radiate as much power as the theoretically maximum absorbed power.

This statement is valid also for systems of different geometrical configurations, where the maximum absorbed power is less than 100 percent of the incident wave power, provided **the required optimum oscillation can be realised**, that is, when no physical amplitude limitation, or other constraint, prevents the desired radiated wave from being realised.







$$m\ddot{x} + m_{r} * \ddot{x} + R_{r} * \dot{x} + R_{u} \dot{x} + R_{f} \dot{x} + Sx = F_{e}$$
Fe is the wave excitation force (that is, the wave fore which the body experiences if it is not moving)
Useful power
$$\overline{P_{u}} = R_{u} \frac{M^{2}}{M^{2}} = R_{u} \frac{1}{2} \frac{M^{2}}{M^{2}} \qquad M = \dot{x}$$

$$R_{u} \dot{x} = F_{e} - R_{u} \dot{x} - R_{f} \dot{x} - (m + m_{u}) \ddot{x} - Sx$$

$$R_{u} \dot{x}^{2} = F_{e} \dot{x} - R_{u} \dot{x}^{2} - R_{o} \dot{x}^{2} - (m + m_{u}) \ddot{x} \dot{x} - Sx \dot{x}$$

$$- \frac{m + m_{u}}{2} \frac{1}{dt} \dot{x}^{2} - \frac{5}{2} \frac{1}{dt} x^{2}$$

$$= -\frac{d}{dt} (M_{kinethe} + M_{polyntrial})$$

$$\begin{aligned} \overline{P}_{u} = \mathcal{R}_{u} \overline{x}^{\flat} &= \overline{F_{e} x} - \mathcal{R}_{h} \overline{x}^{\flat} - \mathcal{R}_{f} \overline{x}^{\flat} \\ F_{e} \dot{x} &= F_{e} \mu = F_{o} \cos(\omega t) \mathcal{U}_{o} \cos(\omega t - \varphi) \\ & \cos \alpha_{*} \cos \alpha_{2} = \frac{1}{2} \cos(\alpha_{4} - \alpha_{*}) + \frac{1}{2} \cos(\alpha_{*} + \alpha_{*}) \\ F_{e} \dot{x} &= F_{o} \mathcal{U}_{o} \left(\frac{1}{2} \cos \varphi + \frac{1}{2} \cos(2\omega t - \varphi)\right) \\ \overline{F_{e} \dot{x}} &= \frac{1}{2} \overline{F_{o}} \mathcal{U}_{o} \cos \varphi \\ \overline{R_{p} \dot{x}^{\bullet}} &= \frac{1}{2} \mathcal{R}_{f} \mathcal{U}_{o}^{\flat} \\ \overline{P_{u}} &= \frac{1}{2} \overline{F_{o}} \mathcal{U}_{o} \cos \varphi - \frac{1}{2} \mathcal{R}_{n} \mathcal{U}_{o}^{2} - \frac{1}{2} \mathcal{R}_{f} \mathcal{U}_{o}^{2} \end{aligned}$$

In this simple example, at optimum radiated-wave generation, the maximum absorbed energy equals 100 percent of the incident wave energy. Note also that the required, optimum, radiated wave has the same amplitude as the incident wave. Thus,



Observe that, in order to absorb, from the sea, the theoretically maximum wave power, it is necessary that the waveabsorbing oscillating system, at optimum, has an *ability* to radiate as much power as the theoretically maximum absorbed power.

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$$\overline{P_{u}} = \frac{4}{2} F_{0} W_{0} \cos \varphi - \frac{1}{2} R_{u} W_{0}^{2} - \frac{1}{2} R_{y} U_{0}^{2}$$
How does $\overline{P_{u}}$ depend on oraillakin amplitude M_{0} ?
$$\overline{R_{u}}$$

$$\frac{\frac{4}{2} F_{0} W_{0} \cos \varphi}{\frac{1}{2} F_{0} W_{0} \cos \varphi}$$

$$\frac{\frac{1}{2} F_{0} W_{0} \cos \varphi}{\frac{1}{2} R_{u} W_{0}^{2} = hadisked power$$

$$\frac{\frac{1}{2} R_{u} W_{0}^{2} = hadisked power$$

$$\frac{1}{2} R_{y} W_{0}^{2} = hadiske power$$

$$\frac$$



The wave-power "island"

illustrates the real-valued absorbed wave power P_a versus a complex oscillation amplitude U, where $|U|^2 = U U^*$ equals the radiated power P_r . The phase of U is chosen in order to make U real and positive when it has the same phase as the excitation force from the incident wave. The optimum value U_0 is a positive real quantity.

$$P_{a,MAX} = P_{r,OPT} = |U_0|^2$$

 $P_{a,MAX} - P_a = |U_0 - U|^2$

These simple equations are applicable to many different types of wave-energy converters (WECs). Assuming that the power take-off (PTO) machinery is equipped with sufficient control, we may consider *U* to be an independent complex variable. The optimum value U_0 is, however, proportional to the incident wave amplitude *A*, and, moreover, it is a function of β , the angle of wave incidence.

Three wave-power inventor pioneers



Yoshio Masuda (1925 – 2009) Started already in 1947 with experiments to test devices for utilising wave energy in Japan.



Stephen Salter (1938 –) started 1973 wave-power research at the University of Edinburgh, Scotland.



Kjell Budal (1933 – 1989) initiated in 1973 wave-power research at NTH (part of pre-NTNU university), Trondheim, Norway.























Phase-controlled power-buoy model (type E) under test in the Trondheim Fjord, 1983. Video clip [also on <u>http://folk.ntnu.no/falnes/w_e/.]</u>







Photo: NORWAVE AS, 1986

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OED (Ministry of Petroleum and Energy) issued 1987 two reports on NORWAVE's and Kvaerner's wave-power prototypes, 40 km off Bergen. One report, "*Norwegian wave power plants 1987*", with text in Norwegian and English, was open.

The other report, "Bølgekraftverk Toftestallen: Prosjektkomiteens sluttrapport 31.12.1987", had only closed distribution. It contained more detailed information, in the Norwegian language, only.



NORSKE BØLGEKRAFTVERK

1987 OED MPE



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The other report, "Bølgekraftverk Toftestallen: Prosjektkomiteens sluttrapport 31.12.1987", had only closed distribution. It contained more detailed information, in the Norwegian language, only. By end of 1988 Kværner's 500 kW OWC prototype had delivered 29 MWh to the local utility Nordhordland Kraftlag.

It seems that the installed power capacity was much too large!





By end of 1991 NORWAVE's 350 kW TapChan prototype had delivered 691 MWh to the local utility Nordhordland Kraftlag.

Energy deliveries as informed by Nordhordland Kraftlag in letter 1993

In the early 1980s Kværner Brug AS planned a multiresonant OWC WEC standing on 25 m deep sea bed.

Figure from Stortingsmelding [White Paper] nr. 65 (1981-82): Om nye fornybare energikilder i Norge [On new renewable energy sources in Norway].

Fig. 7. Kværners svingende vannsøyle er bygget i betong. Vannsøylen inne i konstruksjonen settes i bevegelse (pil) og driver en luft-turbin bakerst til høyre.

Kværner Brug's 500 kW WEC of the OWC type in a very steep cliff on island Toftøy, 40 km NW from Bergen. The red part, below the generator housing, is the housing for a self-rectifying air turbine.

Constructed during 1985 and destroyed by a storm during the last week of 1988.















David Ross, in his 1995 book "Power from the Waves" reports (p.180) from a wave-energy meeting in Brussels 1991:



The discussion saw another round in the debate - - - about whether it was best to go to sea sooner or later. Professor Salter insisted:

I don't want to be the first wave power device at sea. I want to be the last one. I want to make all the mistakes in private, with instruments to tell me what mistakes I have made so that I don't do it again. I want to do all the difficult things in the laboratory. There was enthusiasm for air ships, but the R101 crashed. Airships finally died when the Hindenburg died. If you had a spectacular disaster with one wave energy device, you could drag everything down, too.



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